## Lesson 2, Diffusion processes

## 1 Introduction

A diffusion process is a kind of random process. We let $X_{t}$ be the value of the process at time $t$. Most interesting diffusions have more than one component: $X_{t}=\left(X_{1, t}, \ldots, X_{d, t}\right) \in \mathbb{R}^{d}$. You can picture $X_{t}$ as the location of a particle moving at random, but not completely randomly. The motion is governed by a deterministic component called drift or infinitesimal mean, and by a random noise. We specify a diffusion process by giving the drift and the noise level, both as a function of $X_{t}$. The model is a stochastic differential equation, or $S D E$. It becomes an ordinary differential equation (ODE) if the noise coefficient is zero. For an Ito type $\mathrm{SDE}, X_{t}$ is a martingale if the drift is zero.

This lesson discusses only the case $d=1$, so $X_{t}$ is just a number. A process is a diffusion (diffusion process) if it has these properties

1. The time variable $t$ is continuous. For convenience we often imagine that the process starts at time $t=0$ and is defined for every (real number) $t>0$.
2. It is a Markov process. This means that $X_{t_{0}}$ determines the distribution of $X_{t}$ for $t>t_{0}$ "completely". The Markov property is explained in more detail below.
3. $X_{t}$ is a continuous function of $t$. There are no "jumps".

Many random processes are modeled as diffusions, either exactly or approximately.

The Markov property has to do with conditional expectation and conditional probability of the future of a path given its present and its past. The path over an interval $\left[t_{1}, t_{2}\right]$ (with $t_{2}>t_{1}$ ) is denoted by $X_{\left[t_{1}, t_{2}\right]}$. At time $t_{1}$, the past is the time interval $\left[0, t_{1}\right]$. The "information" from the past consists of the path $X_{\left[0, t_{1}\right]}$. The information of the present is the single value $X_{t_{1}}$. The Markov property concerns the path in the future of $t_{1}$. It is that the distribution of $X_{\left[t_{1}, t_{2}\right]}$ conditional on the past is the same as the distribution conditional on the present. The distribution conditional on $X_{\left[0, t_{1}\right]}$ is the same as the distribution conditional on $X_{t_{1}}$.

You specify a Markov process by giving its transition probability distributions. Suppose $t_{2}>t_{1}$ and $X_{t}$ is a Markov process. Write $Y$ for $X_{t_{2}}$ and $X$ for $X_{t_{1}}$. The probability density for $X_{t_{2}}$ conditional on $X_{t_{1}}$ is $G\left(y, t_{2}, x, t_{1}\right)$. This is a probability density in the $y$ variable in the sense that

$$
\operatorname{Pr}\left(y \leq X_{t_{2}} \leq y+d y\right)=G\left(y, t_{2}, X_{t_{1}}, t_{1}\right) d y
$$

It may be thought of as the density for transitions

$$
\left(X_{t_{1}}=x \text { at time } t_{1}\right) \longrightarrow\left(Y=X_{t_{2}} \text { at time } t_{2} \text { near } y\right)
$$

In Lesson 1 we used Bayes' rule for Brownian motion to write the joint density for $X_{t}$ at two times. The same reasoning applies here, except that the transition density may depend on $t_{1}$ and $t_{2}$ separately, not just on the length of the time interval $t_{2}-t_{1}$. The result is

$$
u_{2}\left(x_{1}, t_{1}, x_{2}, t_{2}\right)=u\left(x_{1}, t_{1}\right) G\left(X_{t_{1}}, t_{2}, X_{t_{1}}, t_{1}\right) .
$$

The joint density of $X_{t_{2}}$ and $X_{t_{1}}$ is equal to the marginal density of $X_{t_{1}}$ multiplied by the conditional density of $X_{t_{2}}$ given $X_{t_{1}}$. For times $t_{1}<t_{2}<t_{3}$, similar reasoning leads to

$$
u_{3}\left(x_{1}, t_{1}, x_{2}, t_{2}, x_{3}, t_{3}\right)=u\left(x_{1}, t_{1}\right) G\left(x_{2}, t_{2}, x_{1}, t_{1}\right) G\left(x_{3}, t_{3}, x_{2}, t_{2}\right)
$$

This is the joint PDF for the three values $X_{1}=X_{t_{1}} . X_{2}=X_{t_{2}}$ and $X_{3}=X_{t_{3}}$. If $X$ is not a Markov process, this last formula is more complicated. The "transition density" for $X_{3}$ would depend on both $X_{2}$ and $X_{1}$.

In continuous time, we need something like the transition distribution for time $d t$. We will see in future lessons that this is related to what is called the generator of the process. For now, it suffices to say that for diffusion processes, the time $d t$ transitions are determined by what is called here the infinitesimal mean and infinitesimal variance. The proper names for these are drift and quadratic variation. The infinitesimal mean is defined by (slightly more formal versions are given below)

$$
\begin{equation*}
\mathrm{E}\left[X_{t+d t}-x \mid X_{t}=x\right]=a(x, t) d t \tag{1}
\end{equation*}
$$

The infinitesimal variance is defined by

$$
\begin{equation*}
\operatorname{var}\left[X_{t+d t} \mid X_{t}=x\right]=v(x, t) d t \tag{2}
\end{equation*}
$$

There is a theorem like the central limit theorem that says that a diffusion process is completely determined by its infinitesimal mean and variance.

Some Markov processes are not diffusions. Just Brownian motion is the simplest diffusion, the Poisson arrival process is the simplest non-diffusion continuous time Markov process. The Poisson arrival process models random events called "arrivals". The number of arrivals from time zero to time $t>0$ is $N_{t}$. The probability of an arrival in $(t, t+d t)$ is $\lambda d t$, with $\lambda$ being the arrival rate parameter. Arrivals in disjoint time intervals are independent. The time to the first arrival has probability density $\lambda e^{-\lambda t}$ (we will see). This arrival process satisfies (1) and (8) with $a(x, t)=\lambda$ and $v(x, t)=\lambda$. The infinitesimal mean and variance do not determine the process completely unless we also know that the process is a diffusion.

A diffusion is a dynamic stochastic model of something. We interpret the Markov property as saying that the model is "complete" in that the state of
the system at time $t$, which is $X_{t}$, contains all the information about the past that is relevant for predicting the future. The SDE that describes the diffusion process is written in Ito form as

$$
\begin{equation*}
d X_{t}=a\left(X_{t}\right) d t+b\left(X_{t}\right) d W_{t} \tag{3}
\end{equation*}
$$

In this SDE, $a(x)$ is the drift and $b(x)$ is the noise coefficient. If $b(x)=0$ (no noise), this might be rewritten in more familiar ODE form as

$$
\begin{equation*}
\frac{d X_{t}}{d t}=a\left(X_{t}\right) \tag{4}
\end{equation*}
$$

More general diffusions are not written in ODE form because $X_{t}$ is not a differentiable function of $t$ even though $X_{t}$ is a continuous function of $t$.

The $d W$ in the noise term has two interpretations. One is as the "differential increment" of Brownian motion. ${ }^{1}$ That is, $d W_{t}=W_{t+d t}-W_{t}$. We saw in Lesson 1 that $d W_{t}$ is a Gaussian random that is independent of everything up to time $t$. The mean is zero and the variance is $d t$. Therefore, the mean of $b\left(X_{t}\right) d W_{t}$ is zero (because $d W_{t}$ is independent of $X_{t}$ ) and the variance is

$$
\operatorname{var}\left(b\left(X_{t}\right) d W_{t} \mid X_{t}\right)=\mathrm{E}\left[b\left(X_{t}\right)^{2}\left(d W_{t}\right)^{2} \mid X_{t}\right]=b^{2}\left(X_{t}\right) d t
$$

The infinitesimal variance of a diffusion is the square of the noise term in the SDE.

The other interpretation does not connect the noise to Brownian motion. In this interpretation, $d W_{t}$ is just a convenient way to write "mean zero, variance $d t$, independent of whatever happened before time $t$ ". The first interpretation (the strong form is helpful in technical analysis of diffusion processes. The second (the weak form), is more useful for modeling. In modeling, you are interested in the process $X_{t}$ on its own, and not in relation to some idealized Brownian motion that is not part of the system you are modeling. In the strong form, $X$ is a function of $W$. In the weak form, $X$ lives on its own.

Brownian motion itself is the simplest interesting diffusion. It was called $X_{t}$ in Lesson 1, but it is often called $W_{t}$ to distinguish it from other diffusion processes. Recall that standard Brownian motion is Gaussian with the properties

1. $W_{0}=0$.
2. If $t_{2}>t_{1}$, then $\mathrm{E}\left[W_{t_{2}} \mid W_{t_{1}}\right]=W_{t_{1}}$ (the martingale property).
3. $\operatorname{var}\left[W_{t_{2}}-W_{t_{1}}\right]=t_{2}-t_{1}$.

Brownian motion is a particularly simple Markov process in which the increment is independent of the present and the past. The increment between time $t_{1}$ and time $t$ (with $t>t_{0}$ ) is $\Delta W=W_{t}-W_{t_{1}}$. The increment of Brownian motion is independent of $W_{\left[0, t_{1}\right]}$. The SDE that describes Brownian motion has zero drift ( $a=0$ ) and constant noise coefficient $b=1$.

[^0]The Ornstein Uhlenbeck process has a linear drift that seeks to return $X$ to zero and a constant noise:

$$
\begin{equation*}
d X_{t}=-\gamma X_{t} d t+\sigma d W_{t} \tag{5}
\end{equation*}
$$

The mean reversion rate coefficient $\gamma$, and the noise coefficient $\sigma$ are constants of the model. The OU process is often a good model of a system with a stable equilibrium that is subject to small outside disturbances. It was used by Einstein as a model of the velocity of a small particle in a fluid, with $X_{t}$ being the velocity at time $t$. A moving particle will slow because of friction with the fluid, which is modeled by $-\gamma X$, the friction force is proportional to the velocity. The particle also is subject to random forces caused by collisions from water molecules. In Einstein's simple model the amount of noise is constant, independent of time and the speed of the particle. This is $\sigma d W$.

Geometric Brownian motion models exponential growth (or decay) in the presence of noise. It differs from the OU process in that the noise is proportional to the level. It is defined by a growth rate parameter $\mu$ and a volitility parameter $\sigma$.

$$
\begin{equation*}
d S_{t}=\mu S_{t} d t+\sigma S_{t} d W_{t} \tag{6}
\end{equation*}
$$

Geometric Brownian motion is a simple model of the random price of a share of stock through time. If there is no noise, then the stock is a simple exponential. The noise is made proportional to the level so that the value of $N$ "shared" of stock doesn't change under a "stock split". A stock split is replacing each share by two shares worth half the amount: $n \rightarrow 2 n$, and $S \rightarrow \frac{1}{2} S$. This is our first model with multiplicative noise, which means that $b(x)$ is not a constant but varies with $x$. The generic $X_{t}$ is replaced by $S_{t}$ (for "stock") for geometric Brownian motion.

This Lesson begins the discussion of diffusion processes. The next section gives a more technical definition of the Markov property, drift and noise. The main goal of this lesson is the partial differential equations (PDEs) related to $X_{t}$, which are the backward equation and the forward equation.

## 2 Diffusions

It is a theorem (not proved in this course) that a diffusion process is determined by the infinitesimal mean and infinitesimal variance. Infinitesimal mean is often called drift and infinitesimal variance is called quadratic variation. "Information about the past" of $t$ is denoted $\mathcal{F}_{t}$. The precise definition of $\mathcal{F}_{t}$ is not important yet. The important thing here is that if $A$ is anything determined by the path, then

$$
\begin{equation*}
\mathrm{E}\left[A \mid X_{\left[0, t_{1}\right]}\right]=\mathrm{E}\left[A \mid \mathcal{F}_{t_{1}}\right] \tag{7}
\end{equation*}
$$

Suppose $d t>0$ is an infinitesimal increment of time. ${ }^{2}$ The corresponding increment of the diffusion process is $d X_{t}=X_{t+d t}-X_{t}$. The infinitesimal mean

[^1]and variance defined above may be written
$$
E\left[d X_{t} \mid \mathcal{F}_{t}\right]=a\left(X_{t}, t\right) d t, \quad E\left[\left(d X_{t}\right)^{2} \mid \mathcal{F}_{t}\right]=v\left(X_{t}, t\right) d t
$$

The infinitesimal variance, $v(x, t)$ is defined by

$$
\begin{equation*}
\operatorname{var}\left[d X_{t} \mid \mathcal{F}_{t}\right]=v\left(X_{t}, t\right) d t \tag{8}
\end{equation*}
$$

For more careful (but still not completely rigorous) mathematical work, we define $\Delta t>0$ to be a small but not infinitely small increment of time and $\Delta X_{t}=X_{t+\Delta t}-X_{t}$ the corresponding small increment of $X$. We will derive simple formal formulas involving differentials using more complicated formulas involving $\Delta X$ and $\Delta t$. For example, much of ordinary (deterministic) calculus may be summarized by saying $(d t)^{2}=0$ even though $d t>0$. For example,

$$
\left.d\left(t^{2}\right)=(t+d t)^{2}-t^{2}=t^{2}+2 t d t+d t^{2}-t^{2}=2 t d t \quad \text { (because } d t^{2}=0\right)
$$

We might then divide by $d t$ to get

$$
\frac{d}{d t} t^{2}=2 t
$$

Here's the same thing done less informally with $\Delta t$.

$$
\Delta\left(t^{2}\right)=(t+\Delta t)^{2}-t^{2}=2 t \Delta t+O\left(\Delta t^{2}\right)
$$

Therefore

$$
\begin{aligned}
\frac{d}{d t} t^{2} & =\lim _{\Delta t \rightarrow 0} \frac{\Delta\left(t^{2}\right)}{\Delta t} \\
& =\lim _{\Delta t \rightarrow 0}\left[\frac{2 t \Delta t+O\left(\Delta t^{2}\right)}{\Delta t}\right] \\
& =\lim _{\Delta t \rightarrow 0}[2 t+O(\Delta t)] \\
& =2 t
\end{aligned}
$$

The definition of $Q=O\left(\Delta t^{p}\right)$ is: there is an $\epsilon>0$ and a $C>0$ so that if $\Delta t \leq \epsilon$ then $|Q|<C \Delta t^{p}$. We say that the quantity $Q$ is "on the order of" $\Delta t^{p}$. We used two facts about "big Oh" orders. One is that if $Q=O\left(\Delta t^{p}\right)$ and $p>1$, then $Q / \Delta t=O\left(\Delta t^{p-1}\right)$. The other is that if $R=O\left(\Delta t^{p^{\prime}}\right)$, with $p^{\prime}>0$ then $R \rightarrow 0$ as $\Delta t \rightarrow 0$. We write " $O\left(\Delta t^{p}\right)$ " instead of " $Q$, with $Q=O\left(\Delta t^{p}\right)$ ". If you did not study mathematical analysis in college these definitions may be confusing at first. But they do not get more complicated than this and should quickly seem natural. You basically treat $O\left(\Delta t^{p}\right)$ as though it were $\Delta t^{p}$,

[^2]even though the truth is more complicated. In the above calculation, we used $O\left(\Delta t^{2}\right) / \Delta t=O(\Delta t)$ and then $O(\Delta t) \rightarrow 0$ as $\Delta t \rightarrow 0$.

In the $\Delta t$ framework, infinitesimal mean is

$$
\begin{equation*}
\mathrm{E}\left[\Delta X_{t} \mid \mathcal{F}_{t}\right]=a\left(X_{t}, t\right) \Delta t+O\left(\Delta t^{2}\right) \tag{9}
\end{equation*}
$$

the infinitesimal variance is

$$
\begin{equation*}
\operatorname{var}\left(\Delta X_{t} \mid \mathcal{F}_{t}\right)=v\left(X_{t}, t\right) \Delta t+O\left(\Delta t^{2}\right) \tag{10}
\end{equation*}
$$

There is a simpler expression

$$
\begin{equation*}
\mathrm{E}\left[\left(\Delta X_{t}\right)^{2} \mid \mathcal{F}_{t}\right]=v\left(X_{t}, t\right) \Delta t+O\left(\Delta t^{2}\right) \tag{11}
\end{equation*}
$$

The formula (10) is equivalent to (11), because, if we assume 11), then

$$
\begin{aligned}
\operatorname{var}\left(\Delta X_{t} \mid \mathcal{F}_{t}\right) & =\mathrm{E}\left[\left(\Delta X_{t}\right)^{2} \mid \mathcal{F}_{t}\right]-\mathrm{E}\left[\Delta X_{t} \mid \mathcal{F}_{t}\right]^{2} \\
& =\mathrm{E}\left[\left(\Delta X_{t}\right)^{2} \mid \mathcal{F}_{t}\right]-\left[a\left(X_{t}, t\right) \Delta t+O\left(\Delta t^{2}\right)\right]^{2} \\
& =v\left(X_{t}, t\right) \Delta t+O\left(\Delta t^{2}\right)
\end{aligned}
$$

Here we use more facts about "big Oh". One is that $\Delta t+O\left(\Delta t^{2}\right)=O(\Delta t)$. Another is that $(O(\Delta t))^{2}=O\left(\Delta t^{2}\right)$. The above calculation should convince you that "big Oh" is a convenient way to reason about small quantities whose exact size isn't relevant.

The formula (11) may be interpreted as saying that $\Delta X_{t}$ is approximately on the order of $\sqrt{\Delta t}$, because $\Delta X^{2}$ is, in the expected value, on the order of $\Delta t$. If $\Delta X$ is on the order of $\Delta t^{\frac{1}{2}}$, then $\Delta X^{4}$ should be on the order of $\Delta t^{2}$. The Poisson arrival process shows that this reasoning is flawed. In fact,

$$
\begin{aligned}
\mathrm{E}\left[\Delta N \mid \mathcal{F}_{t}\right] & =\lambda \Delta t+O\left(\Delta t^{2}\right) \\
\mathrm{E}\left[(\Delta N)^{2} \mid \mathcal{F}_{t}\right] & =\lambda \Delta t+O\left(\Delta t^{2}\right) \\
\mathrm{E}\left[(\Delta N)^{4} \mid \mathcal{F}_{t}\right] & =\lambda \Delta t+O\left(\Delta t^{2}\right)
\end{aligned}
$$

The first two seem fine, but the last one violated the reasoning. More on this in the exercises.

Most of the time you can tell a diffusion process from another Markov process fourth moments. A diffusion process has a fourth moment that follows the informal reasoning:

$$
\begin{equation*}
\mathrm{E}\left[(\Delta X)^{4} \mid \mathcal{F}_{t}\right]=O\left(\Delta t^{2}\right) \tag{12}
\end{equation*}
$$

For most purposes, particularly in this class, you reason about diffusions using the infinitesimal mean (9), the infinitesimal square (11), which is equivalent to the infinitesimal variance (10), and the fourth moment bound (12). Future lessons will have more on the fourth moment bound.

### 2.1 Proofs with big Oh

The statement $Q=O(P)$ as $\Delta t \rightarrow 0$ literally means that $Q$ and $P$ are functions of $\Delta t$ and there is an $\epsilon>0$ and a $C<\infty$ so that $|Q| \leq C P$ if $\Delta t \leq \epsilon$. In a formula like $A=B+O(P)$ as $\Delta t \rightarrow 0$ the $O(P)$ is a quantity $Q$ with $Q=O(P)$. There can be more than one big Oh quantity in a formula, for example $A=B+O(P)+O(R)$ as $\Delta t \rightarrow 0$. Here $O(P)$ represents a quantity $Q_{1}$ with $Q_{1}=O(P)$ and $O(R)$ represents a $Q_{2}$ with $Q_{2}=O(R)$. A typical application has $P$ or $R$ being powers of $\Delta t$. Someone who has understood a class in "Mathematical Analysis" ( $\epsilon$ and $\delta$ proofs) will be able to reason with big Oh. These examples are for people who haven't taken such a class or are rusty.
Example 1. Show that $O\left(\Delta t^{p}\right) / \Delta t=O\left(\Delta t^{p-1}\right)$. A solution: Let $P=$ $O\left(\Delta t^{p}\right)$ be the quantity on the left and $R=P / \Delta t$. From the definition, there is an $\epsilon>0$ and a $C>0$ so that $|P| \leq C \Delta t^{p}$ if $\Delta t \leq \epsilon$. Therefore, if $\Delta t \leq \epsilon$, $R \leq C \Delta t^{p} / \Delta t=C \Delta t^{p-1}$.
Example 2. Show that $O\left(\Delta t^{p_{1}}\right) O\left(\Delta t^{p_{2}}\right)=O\left(\Delta t^{p_{1}+p_{2}}\right)$. A solution: Call the two functions $P_{1}=O\left(\Delta t^{p_{1}}\right)$ and $P_{2}=O\left(\Delta t^{p_{2}}\right)$. By the definitions, there is an $\epsilon_{1}>0$ and a $C_{1}$ so that $\left|P_{1}\right| \leq C_{1} \Delta t^{p_{1}}$ if $\Delta t \leq \epsilon_{1}$. There also is an $\epsilon_{2}$ and $C_{2}$ for $P_{2}$. Take $\epsilon=\min \left(\epsilon_{1}, \epsilon_{2}\right)$. The min of two positive numbers is a positive number. If $\Delta t \leq \epsilon$ then $\Delta t \leq \epsilon_{1}$ and $\Delta t \leq \epsilon_{2}$. Therefore $\left|P_{1}\right| \leq C_{1} \Delta t^{p_{1}}$ and $\left|P_{2}\right| \leq C_{2} \Delta t^{p_{2}}$. Therefore $\left|P_{1}, P_{2}\right| \leq C_{1} C_{2} \Delta t^{p_{1}+p_{2}}$ if $\Delta t \leq \epsilon$. This proves that $\left|P_{1}, P_{2}\right| \leq C \Delta t^{p_{1}+p_{2}}$, with $C=C_{1} C_{2}$. Comment: Example 1 is a special case of Example 2, with $p_{2}=-1$.
Example 3. Is it true that $O\left(\Delta t^{p_{1}}\right) / O\left(\Delta t^{p_{2}}\right)=O\left(\Delta t^{p_{1}-p_{2}}\right)$ ? A solution: It's not true. If $\left|P_{2}\right| \leq C_{2} \Delta t^{p_{2}}$ it might be that $P_{2}$ is much smaller than this. For example, $\Delta t^{2}=O(\Delta t)$. If $P_{2}=\Delta t^{2}$, and $P_{1}=\Delta t^{2}$, then $P_{1} / P_{2}$ is not order $p_{1}-p_{2}=2-1=1$. It is common to write $P=O\left(\Delta t^{p}\right)$ to imply that $P$ is about that size and not very much smaller. But this is not part of the "big Oh" definition.
Example 4. Suppose that $f(x)$ and $f^{\prime}$ and $f^{\prime \prime}$ are continuous functions of $x$. Show that $f(x)=f(0)+x f^{\prime}(0)+O\left(x^{2}\right)$ as $x \rightarrow 0$. Comment: we use $x$ instead of $\Delta t$ as the variable that is going to zero. A solution: One form of the Taylor series remainder formula is

$$
f(x)=f(0)+x f^{\prime}(0)+\frac{1}{2} x^{2} f^{\prime \prime}(\xi), \quad|\xi| \leq x
$$

Here, we know there is a $\xi$ but do not know its value. Since $f^{\prime \prime}$ is continuous, there is a $C$ so that

$$
\max _{|x| \leq 1}\left|f^{\prime \prime}(x)\right|=C<\infty .
$$

Take $\epsilon=1$ and use the $C$ given.

### 2.2 Cauchy Schwarz inequality

Suppose $A$ and $B$ are two random variables. The Cauchy Schwarz inequality is

$$
\begin{equation*}
\mathrm{E}[A B] \leq \sqrt{\mathrm{E}\left[A^{2}\right] \mathrm{E}\left[B^{2}\right]} \tag{13}
\end{equation*}
$$

The proof is a clever trick. For every real number $t$,

$$
\mathrm{E}\left[(A-t B)^{2}\right] \geq 0
$$

If a random quantity is non-negative, then its expected value cannot be negative. Calculating, we get

$$
0 \leq \mathrm{E}\left[A^{2}\right]-2 t \mathrm{E}[A B]+t^{2} \mathrm{E}\left[B^{2}\right]
$$

Since the right side is non-negative for every $t$, its minimum is non-negative. Minimizing over $t$ (differentiate with respect to $t$, set the derivative to zero, solve for $t$ ), the minimum is achieved at

$$
t_{*}=\frac{\mathrm{E}[A B]}{\mathrm{E}\left[B^{2}\right]}
$$

This gives

$$
0 \leq \mathrm{E}\left[A^{2}\right]-\frac{\mathrm{E}[A B]^{2}}{\mathrm{E}\left[B^{2}\right]}
$$

Finally, multiply by the non-negative quantity $\mathrm{E}\left[B^{2}\right]$ and you get

$$
\mathrm{E}[A B]^{2} \leq \mathrm{E}\left[A^{2}\right] \mathrm{E}\left[B^{2}\right]
$$

The square root form of this is (13).
The Cauchy Schwarz inequality implies an inequality involving variance and covariance. Suppose

$$
\bar{A}=\mathrm{E}[A], \bar{B}=\mathrm{E}[B]
$$

Replace $A$ with $A-\bar{A}$ and $B$ with $\bar{B}$. The covariance is

$$
\operatorname{cov}(A, B)=\mathrm{E}[(A-\bar{A})(B-\bar{B})]
$$

The Cauchy Schwarz gives

$$
\operatorname{cov}(A, B) \leq \sqrt{\operatorname{var}(A) \operatorname{var}(B)}
$$

This may be re-written in terms of the correlation coefficient between two random variables:

$$
\operatorname{corr}(A, B)=\frac{\operatorname{cov}(A, B)}{\sqrt{\operatorname{var}(A) \operatorname{var}(B)}}
$$

This is a dimensionless measure of the statistical relationship between $A$ and $B$. The Cauchy Schwarz inequality implies that

$$
-1 \leq \operatorname{corr}(A, B) \leq 1
$$

Absolute number bounds make sense for correlation because it is dimensionless.
We are interested in the Cauchy Schwarz inequality here because of something technical we are about to do. There will soon be a Taylor series calculation
up to order $\Delta X^{2}$ with an error that is of order $\Delta X^{3}$. We need to "bound" $\Delta X^{3}$ in terms of the second and fourth moments. For this, apply Cauchy Schwarz with $A=|\Delta X|$ and $B=\left|\Delta X^{2}\right|$. The result is

$$
\begin{aligned}
\mathrm{E}\left[|\Delta X|^{3}\right] & =\mathrm{E}\left[|\Delta X|(\Delta X)^{2}\right] \\
& \leq \sqrt{\mathrm{E}\left[(\Delta X)^{2}\right] \mathrm{E}\left[(\Delta X)^{4}\right]}
\end{aligned}
$$

If $X$ is a diffusion process, then we can use the variance bound (11) and the fourth moment bound (12), and some "big Oh calculations" to get

$$
\begin{equation*}
\mathrm{E}\left[|\Delta X|^{3}\right]=\sqrt{O(\Delta t) O\left(\Delta t^{2}\right)}=O\left(\Delta t^{3 / 2}\right) \tag{14}
\end{equation*}
$$

## 3 Backward equation

Suppose $V(x)$ is a payout function. The corresponding value function, $f(x, t)$, is defined for $t \leq T$, by

$$
\begin{equation*}
f(x, t)=\mathrm{E}\left[V\left(X_{T}\right) \mid X_{t}=x\right]=\mathrm{E}_{x, t}\left[V\left(X_{T}\right) \mid X_{t}=x\right] \tag{15}
\end{equation*}
$$

The terms "payout" and "value" come from financial applications, but the ideas are more general than finance. The value function $f$ depends on the payout function $V$ and also on the diffusion process $X$. The goal of this section is to show that the value function satisfies the partial differential equation called the backward equation

$$
\begin{equation*}
\partial_{t} f+a(x) \partial_{x} f+\frac{1}{2} v(x) \partial_{x}^{2} f=0 \tag{16}
\end{equation*}
$$

Be careful not to confuse the payout function $V(x)$ with the infinitesimal variance $v(x)$.

The derivation of the backward equation has two steps. The first step is the tower property, also called the law of total probability. This allows us to express the values $f(x, t)$ in terms of the values $f\left(x, t_{1}\right)$ with $t_{1}>t$. The Markov property also enters. It implies that

$$
\begin{equation*}
\mathrm{E}\left[V\left(X_{T}\right) \mid X_{t}=x \text { and } X_{t_{1}}=y\right]=\mathrm{E}\left[V\left(X_{T}\right) \mid X_{t_{1}}=y\right]=f\left(y, t_{1}\right) \tag{17}
\end{equation*}
$$

This leads to the equation

$$
\begin{equation*}
f(x, t)=\mathrm{E}_{x, t}\left[f\left(X_{t_{1}}, t_{1}\right)\right] . \tag{18}
\end{equation*}
$$

The value function at time $t$ is represented as the expected value of the value function at a future time $t_{1}>t$.

The second step is to apply the tower property with $t_{1}=t+\Delta t$ and do Taylor series calculations in $\Delta t$. There is an increment $\Delta X$ corresponding to the time increment $\Delta t$. We will have to expand $f$ to first order in $\Delta t$ and to second order in $\Delta X$. This is because $\mathrm{E}\left[\Delta X^{2}\right]=O(\Delta t)$.

The tower property is that the expected value of the expected value is the expected value. Suppose $(Y, Z)$ is a two dimensional random variable with a joint PDF $u_{2}(y, z)$. Suppose $V(z)$ is a payout function and $g(y)$ is the conditional expectation of $V(Z)$ given that $Y=y$. Suppose $f$ is the unconditional expectation of $V(Z)$. The tower property is

$$
\begin{equation*}
f=\mathrm{E}[g(Y)] . \tag{19}
\end{equation*}
$$

The overall expectation is the expected value of the conditional expectation. this is a "tower" of expectations and conditioning.

Here are the formulas for (19). The marginal probability density for $Y$ is

$$
u_{1}(y)=\int u_{2}(y, z) d z
$$

The conditional density for $Z$ given $Y=y$ is

$$
u(z \mid y)=\frac{u_{2}(y, z)}{u_{1}(y)}
$$

You can check that $u(z, y)$ is a probability density in $z$ for each $y$ by integrating

$$
\int u(z \mid y) d z=\frac{1}{u_{1}(y)} \int u_{2}(y, z) d z=\frac{u_{1}(y)}{u_{1}(y)}=1
$$

The conditional expectation may be written in several ways:

$$
g(y)=\mathrm{E}_{y}[V(Z)]=\mathrm{E}[V(Z) \mid Y=y]=\int V(z) u(z \mid y) d z
$$

The overall expectation is

$$
f=\mathrm{E}[V(Z)]=\iint V(z) u_{2}(y, z) d y d z
$$

The tower property is that this is the expected value of $g$ :

$$
f=\mathrm{E}[g(Y)]=\int g(y) u_{1}(y) d y
$$

We can verify this by substituting some of the above definitions:

$$
\begin{aligned}
\int g(y) u_{1}(y) d y & =\iint V(z) u\left(z_{y}\right) u_{1}(y) d y d z \\
& =\iint V(z) \frac{u_{2}(y, z)}{u_{1}(y)} u_{1}(y) d y d z \\
& =\iint V(z) u_{2}(y, z) d y d z
\end{aligned}
$$

We apply the tower property with $Z=X_{T}$ and $Y=X_{t_{1}}$. The starting value $X_{t}=x$ is fixed throughout. All these calculations assume that (are conditioned
on) $X_{t}=x$. The conditional expectation, called $g$ in the abstract calculations above, is

$$
g(y)=\mathrm{E}\left[V\left(X_{T}\right) \mid X_{t_{1}}=y \text { and } X_{t}=x\right]
$$

The Markov property makes the second condition on the right irrelevant. Since $T$ is in the future of $t_{1}$, which is in the future of $t$, once we know $X_{t_{1}}=y$, the value of $X_{t}$ is irrelevant for expectations involving $X_{T}$. Therefore (this may be the main step in the whole thing)

$$
g(y)=\mathrm{E}\left[V\left(X_{T}\right) \mid X_{t_{1}}=y\right]=f\left(y, t_{1}\right)
$$

The equation (18) follows from this when we substitute the definition $Y=X_{t_{1}}$.
The second step is the Taylor calculations. Take $t_{1}=t+\Delta t$ and $X_{t_{1}}=$ $x+\Delta x$. We expand $f(x+\Delta x, t+\Delta t)$ in Taylor series about $x$ and $t$. Error "estimates" ${ }^{3}$ in Taylor series are usually "the first neglected terms". Example 4 in the big Oh section is like that. It is also true, if you are careful, for functions of more than one variable like $f(x, t)$. If we assume that partial derivatives of $f$ up to third order are continuous, then

$$
\begin{align*}
f(x+\Delta x, t+\Delta t)= & f+\Delta x \partial_{x} f+\frac{1}{2} \Delta x^{2} \partial_{x}^{2} f+\Delta t \partial_{t} f  \tag{20}\\
& +O\left(|\Delta x|^{3}\right)+O(\Delta t|\Delta x|)+O\left(\Delta t^{2}\right) \tag{21}
\end{align*}
$$

The arguments $x, t$ are left out of every term on the right on the top line (20). We write, $f$ for $f(x, t), \partial_{t} f$ for $\partial_{t} f(x, t)$, etc. One of the "first neglected terms" is $\frac{1}{6} \Delta x^{3} \partial_{x}^{3} f$. This is on the order of $|\Delta x|^{3}$. The other lowest order neglected terms involve $\partial_{t}^{2} f$ and $\partial_{t} \partial_{x} f$. They give the other two error contributions on the second line (21).

We put the expansion (20) (21) into the tower property formula (18) with $t_{1}=t+\Delta t$ and $X_{t_{1}}=x+\Delta X$. We take expected values. Terms at time $t$ with argument $x$ come out of the expectation because they are determined (when $\left.X_{t}=x\right)$. The result is

$$
\begin{aligned}
f= & f+\mathrm{E}_{x, t}[\Delta X] \partial_{x} f+\frac{1}{2} \mathrm{E}_{x, t}\left[\Delta X^{2}\right] \partial_{x}^{2} f+\Delta t \partial_{t} f \\
& +\mathrm{E}_{x, t}\left[O\left(|\Delta X|^{3}\right)\right]+\Delta t \mathrm{E}_{x, t}[O(|\Delta X|)]+O\left(\Delta t^{2}\right)
\end{aligned}
$$

We evaluate the expectations on the top line using the infinitesimal mean and variance formulas (9) and (11). The $O\left(\Delta t^{2}\right)$ error terms in those formulas contribute to the $O\left(\Delta t^{2}\right)$ on the second line.

$$
\begin{aligned}
0 & =a(x) \Delta t \partial_{x} f+\frac{1}{2} v(x) \Delta t \partial_{x}^{2} f+\Delta t \partial_{t} f \\
& +\mathrm{E}_{x, t}\left[O\left(|\Delta X|^{3}\right)\right]+\Delta t \mathrm{E}_{x, t}[O(|\Delta X|)]+O\left(\Delta t^{2}\right)
\end{aligned}
$$

[^3]We showed that $\mathrm{E}_{x, t}\left[|\Delta X|^{3}\right]=O\left(\Delta t^{\frac{3}{2}}\right)$. It is an exercise to show that $\mathrm{E}_{x, t}[|\Delta X|]=$ $O\left(\Delta t^{\frac{1}{2}}\right)$. This gives (since $\left.O\left(\Delta t^{\frac{3}{2}}\right)+O\left(\Delta t^{2}\right)=O\left(\Delta t^{\frac{3}{2}}\right)\right)$

$$
0=a(x) \Delta t \partial_{x} f+\frac{1}{2} v(x) \Delta t \partial_{x}^{2} f+\Delta t \partial_{t} f+O\left(\Delta t^{\frac{3}{2}}\right)
$$

Finally, we divide both sides by $\Delta t$ and take the limit $\Delta t \rightarrow 0$. The result is the backward equation (16).


[^0]:    ${ }^{1}$ Brownian motion is also called the Wiener process, after MIT mathematician Norbert Wiener. The notation $W_{t}$ is for Wiener.

[^1]:    ${ }^{2}$ Mathematicians don't like $d t$ because $d t$ is supposed to be smaller than any positive number and yet not zero. As a mathematical fact, if $Q \geq 0$ and $Q$ is less than any posi-

[^2]:    tive number, then $Q=0$. Our $d t$ is less formal - very small, positive, yet not zero. The English/Irish philosopher George (Bishop) Berkeley mocked Newton's infinitesimals as the "ghosts of departed quantities". "Departed" means dead; $d t$ has died (gone to zero) and yet lives on (isn't zero).

[^3]:    ${ }^{3}$ An estimate in mathematical proofs is not a guess at how large something is, but an upper bound. Any "big Oh" formula is an "estimate" in this sense.

