

PDE - Lecture 8, 11/4/2014

[start with "Lecture 7 notes", which we didn't get to last week]

Last "elliptic" topic: the "big picture" concerning several approaches to existence, then some detail concerning variational methods. I'll focus mainly on $\Delta u = 0$ in Ω , $u|_{\partial\Omega} = g$ but the story for Neumann bc is $\partial\Omega$ similar (and cases of the form $\Delta u = f$ in Ω are easily reduced to the homop pbm $\Delta \tilde{u} = 0$ if f extends to a smooth fn on all \mathbb{R}^n , by taking $\tilde{u} = u - \Phi * f$ and adjusting the bc accordingly).

Note: it isn't enough to "just use the Green's fn", since our argument for the existence of the Green's fn assumed you can solve $\Delta u = 0$ in Ω , $u|_{\partial\Omega} = g$ for nice enough g .

There are 3 main approaches to existence:

(1) Perron's method rests on the observations that

- if $\Delta w \geq 0$ in Ω ("w is subharmonic") and $w \leq u$ at $\partial\Omega$ then $w \leq u$ in Ω (by max prin, since $\Delta(w-u) \geq 0$).

• same stuff even if w is merely convex and " $\Delta w \geq 0$ " is replaced by

$$(*) \quad w(x) \leq \frac{1}{|\partial B_r(x)|} \int_{|y-x|=r} w(y) \, dA_y$$

(since pt of max prin uses only mean value principle; note by Prob Set 6 #1 that $\Delta w \geq 0$ does imply $(*)$).

So: soln can be characterized as

$u =$ largest conv's fn $w \Rightarrow u$ at $\partial \Omega$ and w is subharmonic (in sense of 2nd bullet)

See eg Fritz John's book 3.4.4 for pt that there is such a function, and that it's harmonic, with $u = g$ at $\partial \Omega$.

Advantage of this approach: it permits very general bc and rather singular domains (though some hypothesis is needed!).

Disadvantage: not well-suited to numerical implementation.

(2) Boundary integral method: solve an integral eqn on $\partial\Omega$; eg to solve $\Delta u = 0$ in Ω with $u = g$ at $\partial\Omega$, we can look for $u: \partial\Omega \rightarrow \mathbb{R}$ at the u_g defined ($\in \mathbb{R}^n$, $n \geq 3$) by

$$u_g(x) = \int_{\partial\Omega} \frac{\partial}{\partial \nu_y} \left(\frac{1}{|x-y|^{n-2}} \right) g(y) dA_y$$

(which is definitely harmonic in Ω) has the desired bc at $\partial\Omega$. This is called the representation of u by a "double layer potential".

See Brenner + Lee 2.8.6-8.7 for a good treatment of this topic.

Advantage: works well numerically (in fact it is the method of choice for unbounded Ω , when discretizing space is impractical)

Disadvantage: method is rather special to Laplace's eqn, and relies heavily on linearity of the pde.

(3) Variational principles, eg

$$(a) \quad \begin{array}{l} \Delta u = f \text{ in } \Omega \\ u = g \text{ at } \partial\Omega \end{array} \iff \min_{\substack{u=g \\ \text{at } \partial\Omega}} \int_{\Omega} \frac{1}{2} |\nabla u|^2 + f u \, dx$$

$$(b) \quad \begin{array}{l} \Delta u = f \text{ in } \Omega \\ \frac{\partial u}{\partial n} = g \text{ at } \partial\Omega \end{array} \iff \min_u \int_{\Omega} \frac{1}{2} |\nabla u|^2 + f u \, dx - \int_{\partial\Omega} u g \, dA.$$

[assuming consistency]

Advantages: easy to implement numerically (by minimizing over a finite-dim'l class of u 's), and extends easily to many nonlinear problems.

Disadvantages: not good for exterior domains, tho: not every pde comes from a var'l pbn. [But: for lin eqns there's a fix - known as the Lax-Milgram Lemma - that permits even eqns not assoc to var'l pbns to be reduced to a linear alg pbn similar to the one assoc to a var'l pbn].

Let's go a bit further on var'l principles:

- Fact that the min is achieved is nontrivial; for the Neumann problem (b) for example it requires that data be consistent.

- Convexity w.r.t. ∇u is also important here; for example

$$\min_u \int_0^1 (u_x^2 - 1)^2 + u^2 dx$$

does not achieve its min (the min is 0, but no function can have $u^2=0$ and $(u_x^2-1)^2=0$).

- if $W(\nabla u)$ is a ^{strictly} convex fn of ∇u , then it's easy to see that a sufficiently smooth soln of

$$\operatorname{div} \left(\frac{\partial W}{\partial \nabla u} \right) = 0 \quad \text{in } \Omega, \quad u = g \quad \text{at } \partial\Omega$$

achieves $\min_{\substack{u=g \\ \text{at } \partial\Omega}} \int_{\Omega} W(\nabla u) dx$

(we'll have a HW prob on this; note that for $W(\nabla u) = |\nabla u|^2$ it was on Problem Set 5).

Rigorous treatment of existence by this method is PDE II material (but you can find it in Chapter 11 of Greenlee + Lee, for example).

Best lets discuss why variational principles like (a) + (b) above lead to successful numerical schemes, focusing for simplicity on

$$\Delta u = f \text{ in } \Omega, \quad u = 0 \text{ at } \partial\Omega.$$

Numerically, we could

• minimize $\int_{\Omega} \frac{1}{2} |\nabla u|^2 + fu$ over some

finite-dimensional space \mathcal{S} of functions on Ω that all vanish at $\partial\Omega$,

The condition of optimality is that opt'l $u_{\mathcal{S}}$ satisfies

$$(\forall) \quad u \in \mathcal{S} \quad \text{and} \quad \int_{\Omega} \langle \nabla u, \nabla \varphi \rangle + f\varphi = 0 \quad \text{for all} \quad \varphi \in \mathcal{S}$$

which amounts to a finite-dimensional linear algebra problem.

Typical choices of \mathcal{S}

- span of 10^4 bases φ_n from a well-chosen basis (eg eigenvectors of the Laplacian, if known, eg on a rectangle; or perhaps a wavelet basis, down to a specified

level of fine-ness)

or

- piecewise linear h_n s on a specified triangulation (this is the simplest example of a finite element method). Note that a piecewise linear h_n is entirely determined by its nodal values; also, there are no restrictions on the nodal values (a piecewise linear h_n is automatically continuous across edges)

How well does this work? Most basic result is that its performance is limited only by how well the true soln can be approximated in S ; in fact if u_S solves $(*)$ + u solves pde then

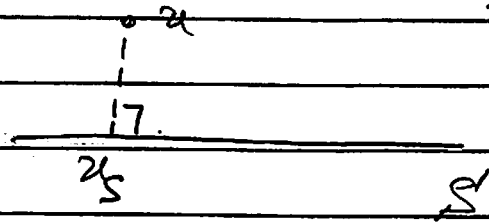
$$\int_{\Omega} |\nabla(u - u_S)|^2 = \min_{v \in S} \int_{\Omega} |\nabla(u - v)|^2$$

Proof of this: From 1st order optimality condns for the variational principle + its restrn to S , we get

$$\int_{\Omega} \langle \nabla(u_S - u), \nabla \varphi \rangle = 0 \quad \text{for all } \varphi \in S$$

This means $u_S =$ orthogonal projn of u onto S w.r.to

The inner product $(u, v) = \int_{\Omega} (\nabla u, \nabla v)$.



Thus, from linear algebra, u_S is the closest pt in S to u . Done!

[Here's a direct version of the argument:
for any $v \in S$ write $u - v = (u - u_S) + (u_S - v)$. Then

$$\int_{\Omega} \frac{1}{2} |\nabla(u-v)|^2 = \int_{\Omega} \frac{1}{2} |\nabla(u-u_S)|^2 + \int_{\Omega} \frac{1}{2} |\nabla(u_S-v)|^2 + \int \nabla(u-u_S), \nabla(u_S-v) \rangle.$$

Last term vanishes since $u_S - v \in S'$. So

$$\int_{\Omega} |\nabla(u-v)|^2 \geq \int_{\Omega} |\nabla(u-u_S)|^2.$$

with equality only when $\int_{\Omega} |\nabla(u_S-v)|^2 = 0$, i.e. $u_S = v$ (using the bc).

For more on this topic see 3.8.5 of Strauss and 2.11.5 of Guenther + Lee (for short discussions) and 1st 50 pp of G. Strang + G. Fix, "An analysis of the finite element method" (a very readable tutorial with much more detail.)