Stochastic Calculus, Courant Institute, Fall 2012
http://www.math.nyu.edu/faculty/goodman/teaching/StochCalc2012/index.html
Always check the class message board on the blackboard site from home.nyu.edu before doing any work on the assignment.

## Assignment 6, due November 12

Corrections: (none yet.)

1. Suppose $X_{t}=\left(X_{1, t}, X_{2, t}, \ldots, X_{n, t}\right)$, where the $X_{k, t}$ are independent standard Brownian motions. The distance of $X_{t}$ from the origin is

$$
R_{t}=\left(X_{1, t}^{2}+\cdots+X_{n, t}^{2}\right)^{1 / 2}
$$

We calculate dynamics of the Bessel process $R_{t}$. This may be found in Wikipedia, but please try to do the problem independently.
(a) Suppose $f\left(x_{1}, \ldots, x_{n}\right)$ is some smooth function, and the $X_{k, t}$ are independent standard Brownian motions. Derive Ito's lemma

$$
d f_{t}=\sum_{k=1}^{n} \partial_{x_{j}} f d X_{j}+\frac{1}{2} \sum_{j=1}^{n} \partial_{x_{j}}^{2} f d t
$$

This is written in vector notation as

$$
d f_{t}=\nabla f \cdot d X_{t}+\frac{1}{2} \triangle f d t
$$

The Laplace operator, or Laplacian, acting on $f$ is

$$
\triangle f=\sum_{j=1}^{n} \partial_{x_{j}}^{2} f
$$

In particular, $\mathrm{E}\left[d f_{t} \mid \mathcal{F}_{t}\right]=\frac{1}{2} \triangle f\left(X_{t}\right) d t$.
(b) For $f(x)=|x|=\left(x_{1}^{2}+\cdots+x_{n}^{2}\right)^{1 / 2}$, calculate $\nabla f$ and $\triangle f$. Hint: $\nabla f$ points directly away from the origin (why?). For $\partial_{x_{j}}^{2} f$, you have to use the chain rule, twice.
(c) For $f(x)=|x|$, calculate $\mathrm{E}\left[d R_{t} \mid \mathcal{F}_{t}\right]=a\left(R_{t}\right) d t$ and $\mathrm{E}\left[d R_{t}^{2} \mid \mathcal{F}_{t}\right]=$ $\mu\left(R_{t}\right) d t$.
2. (Forward equation) Suppose $X_{t}$ is a diffusion process and $u(x, t)$ is the probability density of $X_{t}$ (as a function of $x$ ). Then $u$ satisfies a forward equation, which is a PDE. Early in the course we saw that if $X_{t}$ is Brownian motion, then $u$ satisfies the heat equation

$$
\partial_{t} u=\frac{1}{2} \partial_{x}^{2} u
$$

This exercise suggests how to find the forward equation for more general diffusions. The full derivation by this method is time consuming, so we will do a special case, the case of additive noise. Additive noise means that the coefficient $\mu(x)$ in the infinitesimal variance (equation (2) of Week 7) is independent of $x$.
The first step is to approximate the process $X_{t}$ by a discrete time process $X_{j}^{\Delta t}$. We want $X_{j}^{\Delta t} \approx X_{t_{j}}$, where $t_{j}=j \Delta t$ as always. The approximation will take the form

$$
\begin{equation*}
X_{j+1}^{\Delta t}=X_{j}^{\Delta t}+a\left(X_{j}^{\Delta t}\right) \Delta t+b Z_{j} \tag{1}
\end{equation*}
$$

where the $Z_{j}$ are i.i.d. $Z_{j} \sim \mathcal{N}(0,1)$, and the coefficient $b$ is yet to be determined. We want the approximation to have the properties that

$$
\begin{equation*}
\mathrm{E}\left[\Delta X_{j}^{\Delta t} \mid \mathcal{F}_{j}\right]=a\left(X_{j}^{\Delta t}\right) \Delta t \tag{2}
\end{equation*}
$$

and

$$
\begin{equation*}
\operatorname{var}\left(\Delta X_{j}^{\Delta t} \mid \mathcal{F}_{j}\right)=\mu \Delta t \tag{3}
\end{equation*}
$$

Here, $\mu$ is the constant value of $\mu(x)$ above. The formula (1) satisfies the drift condition (2) automatically. Computing from (1) gives var $\left(\Delta X_{j}^{\Delta t} \mid \mathcal{F}_{j}\right)=$ $b^{2}$. This gives (3) if $b=\sqrt{\mu \Delta t}$. Therefore, our approximation is

$$
\begin{equation*}
X_{j+1}^{\Delta t}=X_{j}^{\Delta t}+a\left(X_{j}^{\Delta t}\right) \Delta t+\sqrt{\mu \Delta t} Z_{j} \tag{4}
\end{equation*}
$$

This approximation satisfies the continuity condition

$$
\mathrm{E}\left[\left(\Delta X_{j}^{\Delta t}\right)^{4} \mid \mathcal{F}_{j}\right] \leq C \Delta t^{2}
$$

It is possible to prove that approximations with these properties converge to the exact stochastic process $X_{t}$ in the sense of distributions.
Let $u_{j}(x)$ be the probability density of $X_{j}^{\Delta t}$. This should be an approximation of the probability density of $X_{t_{j}}$, which is

$$
u_{j}(x) \approx u\left(x, t_{j}\right)
$$

The forward equation PDE that $u(x, t)$ satisfies has the form

$$
\begin{equation*}
\partial_{t} u=L^{*} u \tag{5}
\end{equation*}
$$

where $L^{*}$ is a differential operator in the $x$ variable. We will find $L^{*}$ by finding a formula

$$
\begin{equation*}
u_{j+1}(x)=u_{j}(x)+\Delta t L^{*} u_{j}(x)+O\left(\Delta t^{3 / 2}\right) \tag{6}
\end{equation*}
$$

The technique will be to derive an integral formula for $u_{j+1}$ in terms of $u_{j}$, and then to derive (6) from the integral formula by approximating the integral.
(a) Let $v_{j}(x, y)$ be the joint probability density of $\left(X_{j}^{\Delta t}, X_{j+1}^{\Delta t}\right)$. Here $x$ is the $X_{j}^{\Delta t}$ variable and $y$ is the $X_{j+1}^{\Delta t}$ variable. For example,

$$
\mathrm{P}\left(X_{j}^{\Delta t}>X_{j+1}^{\Delta t}\right)=\int_{-\infty}^{\infty} \int_{y=-\infty}^{x} v_{j}(x, y) d y d x
$$

Write a formula for $v_{j}(x, y)$ in terms of $u_{j}(x)$. This is done using (4) and the formula for a Gaussian density. Use this to find a formula of the form

$$
\begin{equation*}
u_{j+1}(y)=\int_{-\infty}^{\infty} K(x, y, \Delta t) u_{j}(x) d x \tag{7}
\end{equation*}
$$

I with a simple Gaussian explicit formula for $K$.
(b) A "warm up problem" is a problem you know the answer to already that you can use to figure out this new mathematical machinery. In the present case, you know that if $X_{t}$ is Brownian motion, the PDE (5) should be the heat equation. This part of the exercise shows that (6) is satisfied with $L^{*} u=\frac{1}{2} \partial_{x}^{2} u$. In the integral (7), make a Taylor expansion of $u_{j}(x)$ for $x$ near $y$. This may be done in two steps, first to a change of variables $x \leftarrow z=x-y$, then make the Taylor expansion

$$
\begin{aligned}
& u_{j}(y+z)=u_{j}(y)+O(|z|) \quad \text { or } \\
& u_{j}(y+z)=u_{j}(y)+\partial_{x} u_{j}(y) z+O\left(z^{2}\right) \quad \text { or } \\
& u_{j}(y+z)=u_{j}(y)+\partial_{x} u_{j}(y) z+\frac{1}{2} \partial_{x}^{2} u_{j}(y) z^{2}+O\left(|z|^{3}\right) \quad \text { or } \\
& u_{j}(y+z)=u_{j}(y)+\partial_{x} u_{j}(y) z+\frac{1}{2} \partial_{x}^{2} u_{j}(y) z^{2}+\frac{1}{6} \partial_{x}^{3} u_{j}(y) z^{3}+O\left(z^{4}\right)
\end{aligned}
$$

You know you have gone far enough when the remainder term contributes something smaller than $O(\Delta t)$. That is, $\int K(y+z, y, \Delta t)|z|^{p} d z=$ $O\left(\Delta t^{3 / 2}\right)$.
(c) Do the problem now under the hypothesis that $a(x)$ is independent of $x$. Let $x^{*}(y)$ be the $x$ value that maximizes the integrand $K$ over $x$, then let $x=x^{*}(y)+z$. This should give $u_{j+1}(y)=$ $\left.u_{j}\left(x^{*}(y)\right)+\frac{1}{2} \partial_{x}^{2} u_{j}\left(x^{*}\right)\right)$. But $x^{*}(y)=y+O(\Delta t)$, so you can write express $u_{j}\left(x^{*}(y)\right)$ as a Taylor expansion about $u_{j}(y)$ and ignore terms beyond $O(\Delta t)$.
(d) The last step is the hardest. In $K$ you find the expression somewhere $(y-x-a(x) \Delta t)^{2}$. We want to integrate over $x$. We do not need to do the integral exactly, so we can make approximations in this expression if they do not change the answer up to $O(\Delta t)$. Make a Taylor series approximation to $a$ about $y$ ( $y$ is a parameter in the integration) $a(x)=a(y)+a^{\prime}(y)(x-y)+\cdots$ (you figure out how far you need to go). You should get something involving

$$
\frac{u_{j}\left(x^{*}\right)}{1 \pm a^{\prime}(y) \Delta t}+\cdots
$$

This is

$$
u_{j}\left(x^{*}\right)\left(1 \mp a^{\prime}(y) \Delta t\right),
$$

which gives all the contributions of order $\Delta t$. The answer is supposed to be $L^{*} u_{j}(x)=-\partial_{x}\left(a(x) u_{j}(x)\right)+\frac{1}{2} \partial_{x}^{2} u_{j}(x)$.

